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CONTINUATION APPLICATION FOR UNITED STATES PATENT

FOR

Magnetic Etching Process, Especially For Magnetic Or Magnetooptic Recording

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MAGNETIC ETCHING PROCESS, ESPECIALLY FOR
MAGNETIC OR MAGNETOOPTIC RECORDING

The present invention relates to a magnetic
5 etching process.

More particularly, the invention applies
advantageously to ultrahigh-density magnetic recording
(production of discrete magnetic materials, magnetic
memory circuits, magnetically-controllable logic
10 circuits, etc.), optical recording of the read-only
memory type (CDROM, DVDROM, etc.) and production of
magnetically-controllable optical circuits (diffraction
gratings, photonic gap materials, etc.) using a
controlled variation of the optical index component
15 associated with the magnetism.

PRIOR ART

The extraordinary development of multimedia
20 technologies and services in recent years has led to a
race to increase the recording density. In the field of
rewritable disks, although optical (phase change)
technologies are developing rapidly, magnetic
techniques remain the first choice, and most
25 particularly the "hard disk", for its high transfer
rate. However, the current magnetic techniques ought to
be limited to storage densities of $100 \text{ bits}/\mu\text{m}^2$.

One of the limiting factors will especially be
the transition to contact recording, for distances
30 between the read head and the recording medium of less
than 10 nm: there is a trend toward recording
technologies of the "tunnel-effect microscopy"
("STM-like storage") or "near-field" type.

Several technological jumps have been proposed
35 in this direction in recent years, for example near-
field CD-ROM or near-field magnetooptic recording.

In this regard, reference may advantageously be made to the following various publications:

Y. Martin, S. Rishton, H.K. Wickramasinghe, Appl. Phys. Lett. **71**, 1 (1997).

5 Y. Betzig, J.K. Trautman, T.D. Harris, J.S. Weiner, R.L. Kostelak, Science **251**, 1468 (1991).

B.D. Terris, H.J. Mamin, D. Rugar, W.R. Studenmund, G.S. Kino, Appl. Phys. Lett. **65**, 388 (1994).

10 E. Betzig et al., Appl. Phys. Lett. **61**, 142 (1992).

M. Myamoto, J. Ushiyama, S. Hosaka, R. Imura, J. Magn. Soc. Jpn. **19-S1**, 141 (1994).

15 T.J. Silva, S. Schultz, D. Weller, Appl. Phys. Lett. **65**, 658 (1994).

M.W.J. Prinz, R.H.M. Groeneveld, D.L. Abraham, H. van Kempen, H.W. van Kesteren, Applied. Phys. Lett. **66**, 1141 (1995).

Reference may also be made to the publication:

20 B.D. Terris H.J. Mamin, D. Rugar, Appl. Phys. Lett. **68**, 141 (1996) in which it was announced that the company 3M would shortly be commercializing a magnetooptically-read "hard disk" using a solid immersion lens (SIL).

25 However, the main limitation of magnetic techniques should be the "paramagnetic limit", that is to say the size below which the bits will be erased by themselves due to a thermal effect.

30 In the current hard disk technology, the recording medium is a particulate material (magnetic particles in a nonmagnetic matrix, or magnetic particles (grains) separated by nonmagnetic grain boundaries (ME tape)). Now, minimization of the noise necessitates increasing the number of magnetic particles seen by the read head, while these particles must be magnetically decoupled as far as possible. The size of the particles is therefore very much less than the size of a bit. By extrapolating the current data,

the particles would become paramagnetic below 8 nm, thereby limiting the recording density to around 100 bits/ μm^2 .

In magnetooptic recording, the materials used 5 at the present time are amorphous alloys of the rare earth/transition metal type, which could be replaced with Co/Pt multilayers or alloys with the advent of the blue laser. Bits 60 nm in size could actually be written by a thermomagnetic effect in continuous Co/Pt 10 multilayers, but it is probable that noise problems due to the recording medium (domain stability, domain wall roughness) would intervene, at bit sizes very much 15 greater than 60 nm.

To extend this limit, it has recently been 20 proposed to replace the current recording medium materials with discrete materials in which the magnetic bit limits would be geometrically defined by lithographic methods:

either deposition on an etched surface,
25 S. Gadetsky, J.K. Erwin, M. Mansuripur,
J. Appl. Phys **79**, 5687 (1996)

or growth of isolated magnetic particles whose size and position are defined by lithography,

25 S.Y. Chou, M.S. Wei, P.R. Krauss, P. Fischer,
J. Appl. Phys. **76**, 6673 (1994).

The latter technique would allow there to be only a single magnetic particle per bit.

In parallel, pressing techniques based on a matrix defined by electronic lithography have been 30 developed,

S.Y. Chou, P.R. Krauss, P.J. Renstrom, Science **272**, 85 (1996),

Y. Xia, X.M. Zhao, G.M. Whitesides,
Microelecton. Eng. **32**, 255 (1996),

35 which, just as in X-ray or interferential lithography, could in the near future allow mass production of etched media, with patterns very much less than one micron in size over areas of a few cm^2 , probably sufficient for disks of the future.

However, in the current published work, these various techniques have several drawbacks:

1. Whatever the technique adopted, recording in contact mode will require a material having a low and controlled surface roughness: the etched materials proposed up until now will therefore require a final, and probably difficult, planarization step.
2. In the case of near-field magnetooptic recording, sudden variations in optical index (variations in reflectivity) of the etched material will give diffraction effects, which may be manifested by much greater polarization variations than those induced by the magnetic domains - a source of unacceptable noise.
3. A final problem, at very high densities on these etched materials, concerns the following of the track, and it will probably be necessary to develop a specialized "track" for this purpose, but without degrading the points mentioned above.

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PRESENTATION OF THE INVENTION

The subject of the invention is a process for writing on a material, in which said material is irradiated by means of a beam of light ions (that is to say ions having a mass less than 16 units of atomic mass, such as for example He^+ ions), said beam of light ions having an energy of the order of or less than a hundred keV. This process is characterized in that this material comprises a plurality of superposed thin-layers, at least one of the thin layers being magnetic and in that one or more regions having sizes of the order of 1 micrometer or less are irradiated, the irradiation dose being controlled so as to be a few 10^{16} ions/cm² or less, the irradiation modifying the composition of atomic planes in the material at one or more interfaces between two layers of the latter. The magnetic properties of said material, such as, in

particular, its coercivity, its magnetic anisotropy or its Curie temperature, are thus modified.

Typically, a thin layer presents a width of the order of 10 nm or less.

5 The superposed thin layers can advantageously be deposited on a substrate.

They can also be buried in a surface layer.

Such a process allows the aforementioned problems to be solved. In particular:

10 1. The roughness of the original film is unchanged by irradiation and can therefore be adjusted independently. In particular, it may be envisaged to carry out a postirradiation deposition (for the production of devices) under excellent growth 15 conditions (% at an etched surface).

2. The optical index variations remain small for considerable changes in the magnetic properties and can, moreover, be controlled, within a certain range, almost independently of the magnetic variations 20 obtained, by the structure of the substrate or the energy of the ions.

25 3. The effect of the irradiation is cumulative: it is possible to carry out the irradiation several times, and to obtain the same result as in a single time with the cumulative dose. This aspect may be useful when it is desired to irradiate several regions of the specimen with different values, or at 30 different steps in the fabrication of a device.

4. The effect of the irradiation may be easily controlled in real time, by measuring the change in the properties (for example magnetic properties) over a test region.

35 5. The technique is easy to employ for the mass production of recording media, and to do so economically since the tools that it requires to be used are either already used in microelectronics (irradiation) or are under development (lithography by pressing in the case of large areas and of nanometric sizes, for example).

The irradiation may be carried out through a resin mask or with the aid of a focused ion beam.

The aforementioned etching process is advantageously used for the ultrahigh-density magnetic 5 or magnetooptic recording of binary information, and especially for the production of discrete magnetic materials, of magnetic memory circuits or of magnetically-controllable logic circuits.

In particular, the aforementioned process has 10 the advantage of making it possible to write magnetic domains of size very much less than 100 nm and whose position and geometry are perfectly defined and therefore to maximize the signal-to-noise ratio and optimize the track-following problems, while preserving 15 perfectly controlled surface roughness.

In addition, the process proposed by the invention is advantageously used for producing an optical recording of the read-only memory type (CDROM, DVDROM, etc.).

20 It is known in fact that the near-field optical recording techniques will probably have to use smooth writing materials, with a read head flying a few nm above said material (at the present time, 30 nm for a hard disk). Now, the current optical recording 25 techniques of the read-only memory type are not satisfactory: the pressing methods, using dies, may give sizes of less than 100 nm but the recording medium which is obtained is rough; as regards the writing methods using a focused laser beam (ablation, phase 30 change), these do not make it possible to work with bit sizes of the order of or less than 100 nm.

35 Applications other than the recording of binary information may be envisaged. In particular, the magnetic etching process proposed by the invention is advantageously used for the production of magnetically-controllable optical circuits (diffraction gratings, photonic gap materials, etc.) using a controlled variation of the optical index component associated with the magnetism, for the production of sensors (hard

disk read heads, etc.) or magnetic memory circuits (extraordinary Hall-effect memory, magnetoresistive memory, spin-dependent tunnel-effect memory).

In particular, it is known that the emergence
5 of photonic gap materials opens the way to producing
optical devices and that one of the aspects to be
resolved will be that of control of the device. The
process proposed by the invention makes it possible, by
irradiation through a mask, to manufacture a waveguide
10 film made of nonmagnetic material, comprising a regular
array of magnetic units (photonic crystal) having an
optical index which is both slightly different from
that of the host material and magnetically
controllable.

15 In general, the process proposed by the
invention may apply whenever it is advantageous to
define a magnetic element accurately, while maintaining
a very high degree of planarity of the device (for
example, in order to favor subsequent growth).

20 The process proposed by the invention may also
be used for magnetically etching a layer already buried
beneath other, insensitive layers, by adjusting the
irradiation conditions. For example, and by way of
nonlimiting indication, it is possible to produce
25 electrical circuits etched in the same thin-film
magnetic material, and only the important part of which
will remain magnetic, the contact tracks having been
made inactive by irradiation; the coercive field of a
given region of a specimen may be controllably reduced
30 so as to guarantee that the reversal of the
magnetization will always occur under the same
conditions, from the same site.

The process proposed by the invention may *a*
priori be adapted to any material for which a minute
35 variation in the local atomic arrangement can lead to a
large modification in the magnetic properties, that is
to say to transition metal alloys (e.g.: CoPt, NiFe,
etc.), to rare earth/transition metal alloys (e.g.:

TbFeCo, etc.) and to magnetic multilayers (e.g.: Co/Pt, Fe/Tb, etc.), without this list being exhaustive.

Co/Pt multilayers are materials which are potentially of interest for short-wavelength 5 magnetooptic recording in blue light.

DESCRIPTION OF ONE OR MORE EMBODIMENTS

The process of magnetic etching by irradiation 10 is described below in the case of magnetic multilayers irradiated by an ion beam and involves several steps, in which:

- (i) the composition and the roughness at the interfaces and on the surface of the layers are 15 carefully controlled before irradiation;

- (ii) the multilayer structure is irradiated by a light-ion beam, the structural modification induced by the beam being controlled; in particular, the energy density deposited by the beam is controlled 20 by choosing the mass and the energy of the incident ions;

- (iii) the irradiation may be complemented by a suitable thermal processing in order to relax the stresses and/or induce local ordering.

25 In the case of magnetic materials, the effects of the process are important on alloys (transition metal alloys, rare earth alloys and rare earth/transition metal alloys) and on stacks of buried thin layers deposited on a substrate of all types.

30 The process is advantageously employed on Co/Pt multilayers. It should be noted that these materials have already been very widely studied for their properties, firstly their perpendicular magnetic anisotropy and secondly their strong magnetooptic Kerr 35 effect; they therefore constitute advantageous candidates for magnetooptic recording.

In materials based on ultrathin multilayer films, the properties are dominated by the competition between the interface effects and the volume

properties. For example, the easy magnetization direction is given by the sign of an effective anisotropy coefficient K_{eff} which, to a first approximation, is given by:

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$$K_{eff} = -K_d + K_v + \frac{(K_{s1} + K_{s2})}{t_{co}}$$

The first term represents the dipole shape anisotropy ($K_d > 0$), the second term represents the 10 volume anisotropy ($K_v > 0$ in the case of Co) and the last term is due to the interfaces ($K_s > 0$ in the case of the Co/Pt interface), the influence of which varies inversely with the Co thickness t_{co} (K_{s1} and K_{s2} denoting the 15 magnetic anisotropy coefficients of the two interfaces of the Co film. Depending on the sign of K_{eff} , the easy magnetization axis is either the axis perpendicular to the plane of the layers ($K_{eff} > 0$) or the plane of the film. The perpendicular configuration is necessary for magnetooptic recording and will 20 probably become the standard for ultrahigh-density magnetic recording, all techniques included.

The process is limited to irradiation resulting in low energy deposition (small number of atomic displacements at the interfaces that we are interested 25 in). This may be achieved, for example, by light ions (e.g. He^+) of low energy (from a few keV to about a hundred keV). The irradiation firstly modifies the composition of the interface between two layers of material and therefore, in particular, the anisotropy. 30 For the thinnest films (1 or 2 atomic planes) or for higher doses, the composition of the film and hence its volume magnetism are also modified (by transferring atoms from one layer to another): in the particular case of Co/Pt, the Curie temperature of the CoPt alloy 35 decreases with Pt concentration, and becomes below room temperature at around 75% Pt.

For example, the inventors have rendered specimens, having a thickness t_{co} of 0.5 nm,

paramagnetic at ordinary temperature, in a controlled manner, by irradiating, at a (very low) dose of 10^{16} ions/cm², with 30 keV He⁺.

5 The effects of the irradiation were firstly characterized on simple Pt(3.4 nm)/Co(t_{Co})/Pt(6.5 nm)/amorphous substrate (Herasil polished silica, SiO₂/Si, Si₃N₄/Si) sandwiches deposited by sputtering.

10 With the deposition technique used, magnetic films with a perpendicular easy magnetization axis and a perfectly square polar hysteresis cycle (100% remanent magnetization) within the Co thickness range: 0.3 - 1.2 nm are obtained before irradiation.

15 The irradiation of these specimens at He⁺ ion fluences up to around 2×10^{15} atoms/cm², the ions being accelerated to energies of between 5 and 100 keV, makes it possible actually to adjust the magnetic properties of an ultrathin Co layer:

20 1. on 0.5 nm thick layers (approximately 2.25 atomic planes), the main effect is a drop in the Curie temperature, which may fall below room temperature for a dose of the order of 2×10^{16} ions per cm². Below that, the film retains a perpendicular easy magnetization axis and a square loop, but the coercive field of which decreases uniformly when the irradiation 25 dose is increased. Square magnetization loops with coercivities of a few Oe have been obtained. Advantageous applications for the production of low-field sensors may be envisaged;

30 2. on 1 nm thick specimens (approximately 5 atomic planes), the main effect of the irradiation is a tilt of the easy magnetization axis in the plane of the film, combined with a reduction in the interface anisotropy term K_s. The effect is obtained for low doses because the initial thickness is close to that 35 (1.2 nm) at which the tilting effect occurs in the original specimens;

3. on specimens of intermediate thickness (0.8 nm, i.e. 4 atomic planes), the same doses have no visible effect on the hysteresis loop: at these

thicknesses, the Curie temperature is already very high (close to that of bulk Co), and therefore largely insensitive to small modifications of the interface, these thicknesses also being very far from the natural 5 thickness for tilting of the easy magnetization axis. This constitutes a useful characteristic of the process since it makes it possible, on the one hand, to irradiate a bilayer while modifying only one of the layers and, on the other hand, to work at much higher 10 doses, more conducive to homogeneity.

It should be noted that the acceleration energy of the ions has a lesser effect on the modification of the magnetic properties than on the depthwise distribution of the level of displacements in the 15 material. This may allow the process to be employed in thin layers buried at substantially greater depths than those used in the demonstration example.

An essential characteristic of the process proposed is that, although the effect of the 20 irradiation on the magnetism is great, its effect on the optical reflectivity of the specimen remains small.

The contrast is invisible to the naked eye, and barely visible in a good microscope (contrast comparable to that of a domain wall in a Pt/Co/Pt 25 specimen). The smallness of the optical effect is due to the smallness of the induced structural modifications.

Tests on $(\text{Pt}/\text{Co})_6/\text{Pt}$ multilayer stacks were also carried out. The structures of these multilayers 30 (thicknesses, number of Co/Pt periods) were chosen around the values normally used for magnetooptic recording media. Compared with the simple picture of the variation in anisotropy with Co thickness, explained above in the case of the simple films, the 35 effects of the irradiation on the magnetic properties are made more complex in multilayers by the magnetic interaction between the layers, which may be bipolar in origin, or an exchange interaction carried by the conduction electrons in the platinum. The latter

interaction, which is actually manifested by ferromagnetism of the Pt for the interface layers, helps to raise the Curie temperature of the multilayers, especially when the Co thickness is very small. The presence of these two interactions also leads to the existence of quite a wide Co thickness range in which the system is decomposed into regular magnetic domains within which the magnetization is perpendicular ("strip" domain configuration), even for slightly negative K_{eff} values where an easy magnetization plane would be expected.

The tests were carried out on two series of specimens, of the same Co thickness (and therefore the same single layer anisotropy) and the same number of periods, but differing in the thickness of the Pt separating layer:

A series: Pt(2 nm)/[Pt(1.4 nm)/Co(0.3 nm)]₆/Pt(6.5 nm)

B series: Pt(2 nm)/[Pt(0.6 nm)/Co(0.3 nm)]₆/Pt(6.5 nm)

In the case of the B series, the Pt concentration of the alloy after complete interdiffusion would be about 66% (ferromagnetic alloy) while it would be 82% for the A series (nonmagnetic alloy). On the other hand, in the B series, in which the Pt interlayer is thinner, the Co layers are more highly interacting, which in principle makes it easier to obtain the "strip" domain configuration, followed by the easy magnetization plane, by a reduction in the anisotropy.

Over the range of doses tested (up to 10^{16} He/cm² in the case of the A series and 2.6×10^{16} He/cm² in the case of the B series), the irradiation results show qualitatively the same effects for both series: gradual (and easily controllable) transition from a perpendicular easy magnetization axis (with a perfectly square hysteresis loop whose coercive field decreases with the irradiation dose) to a "strip" domain configuration, and then to an easy magnetization plane. As explained above, this tilting takes place at a lower dose for the B series (3×10^{15} He/cm² as opposed

to 6×10^{15} He/cm²). At the doses used, all the specimens remained ferromagnetic at room temperature.

In all the cases described above, no variation in the surface roughness of the specimen could be 5 detected by AFM in air, even for extremely low, of the order of 0.2 nm rms, initial roughnesses.

Tests with irradiation through a resin mask were also carried out.

On Pt(3.4 nm)/Co(0.5 nm)/Pt(6.5 nm)/Herasil 10 simple sandwich specimens, two types of resin were tested:

1. A Shipley negative resin, suitable for submicron lithography by X-ray lithography. The resin had been deposited as a thick (0.8 μm) layer over only 15 half of a specimen and then annealed under the usual conditions. The entire specimen was then irradiated and the resin removed, again under the usual conditions (hot trichloroethylene bath).

The part unprotected by the resin reproduces 20 the effects of the irradiation that were described above, whereas the protected part shows no change in its properties. In principle, using processes already developed elsewhere, the use of the same resin, but with in addition an X-ray lithography step in order to 25 define an array of holes therein, should at the very least make it possible to obtain arrays of magnetically etched bits 0.2 μm in size separated by 0.2 μm , i.e. a recording density of 25 bits per μm^2 , almost 20 times greater than the current densities;

30 2. a PMMA positive resin suitable for electron lithography. The resin was deposited as a layer about 0.85 μm in thickness and in this case was not annealed, something which might have an influence on the quality of the pattern edges. Under the standard annealing 35 conditions for this resin (160°C, 30 min) effects start to appear in the specimens, but annealing of just as good quality is possible at lower temperatures (<120°C), at which the specimens are insensitive). Next, the specimens underwent an electron lithography

step in order to define, as recesses in the resin, an array of lines 1 μm in width, separated by 1 μm , over an area of $800 \times 800 \mu\text{m}^2$. The entire specimen was then irradiated and the resin removed under the standard
5 conditions. Observation in a magnetooptic microscope shows that, at the chosen irradiation dose (10^{16} He/cm^2), the irradiated part becomes paramagnetic at room temperature (this state has the advantage of eliminating the coupling between magnetic regions). The
10 part protected by the resin remains magnetized perpendicularly, with a square loop similar to that of the original specimen.

The same electron lithography process as above was applied to a $\text{Pt}(2 \text{ nm})/\text{[Pt}(0.6 \text{ nm})/\text{Co}(0.3 \text{ nm})]_6/\text{Pt}(6.5 \text{ nm})$ multilayer of the B series in order to create the same array of lines, followed by an irradiation at a dose of $2 \times 10^{15} \text{ He/cm}^2$. However, unlike in the case of the single 0.5 nm Co layer, the two parts (the protected part and the irradiated part)
20 retain a perpendicular magnetization and a square loop with, however, a lower coercive field in the case of the irradiated part. In fact, observation in a magnetooptic microscope clearly shows a reversal of the magnetization in the reverse applied field after
25 saturation, which firstly takes place in the irradiated lines and then propagates into the unirradiated parts (lines and film outside the array). In the intermediate region, magnetic domains artificially created by lithography are therefore obtained. Next, tests were
30 carried out using near-field magnetooptic microscopy, which made it possible to see these artificial domains very precisely. This consequently demonstrates the feasibility of the proposed "contact" recording process. On the other hand, on specimens that were
35 similar but were etched by material ablation, the same near-field microscopy technique reveals only the diffraction effects.

It should be noted that, after irradiation, the PMMA resin becomes more difficult to remove. Residues

remaining along the features introduce roughness and a weak optical contrast of nonmagnetic origin, something which requires an additional stripping procedure in an "oxygen plasma" (a procedure well known in 5 microtechnologies).

Finally, the precision of PMMA-resin electron lithography gives rise to the hope that it will be possible to achieve bit sizes of less than 100 nm, i.e. a density greater than 100 bits/ μm^2 .

10 A series of similar experiments, using masks made of silica and irradiated with He ions under the same conditions as above, has allowed the inventors to expect a resolution (deduced from magnetooptic measurements) of 30 nm on lithographically etched 15 lines.

20 The techniques of the type that have just been described are advantageously used for manufacturing films which include buried magnetic structures, especially for the production of magnetically structured recording media or of magnetoelectronic 25 devices, such as M-RAM memories, logic devices, etc.

They allow planar magnetic etching of buried magnetic layers, which does not modify the surface roughness of the material and makes it possible to 25 control the variations in optical properties, for example to make them negligible.

These techniques can be used for mass production on an industrial scale.

30 Using light ions, which have no etching effect, these can be deeply implanted into the substrate, well below the magnetic layer.

35 The parameter is then the energy deposited per ion along the trajectory - and not the cascades of defects generated by heavy ions - thereby allowing excellent control of the electromagnetic modifications, for high doses, something which gives a homogeneous effect.

Moreover, an easy nucleation region, due to the reversal of the magnetization) and associated with

phenomena occurring at the border of the irradiated region, is intrinsically obtained with the proposed technique. This is a major advantage for controlling and standardizing the magnetization reversal field in 5 an assembly of magnetic "particles", either for a recording medium material or for a memory or logic chip, without limitation.

It should be noted that contrary to heavy ions irradiation techniques in which the atomic 10 modifications are obtained due to the succession chain of defaults created by the heavy ions, the light ions irradiation technique which is herewith presented permits a high control of the magnetic modifications and, for the high irradiation doses, permits to obtain 15 an homogeneous effect.

This is due :

- (i) to the low density of the atomic moves due to the collisions with the atoms at the interface of the thin layers ,
- 20 - (ii) to the low energy transferred during these collisions.